VLT near-infrared spectra of hard serendipitous Chandra sources

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ABSTRACT
We present near-infrared long-slit spectra of eight optically-dim X-ray sources obtained with ISAAC on the Very Large Telescope. Six of the sources have hard X-ray emission with a significant fraction of the counts emerging above 2 keV. All were discovered serendipitously in the fields of three nearby galaxy clusters observed with Chandra, and identified through near-infrared imaging. The X-ray fluxes lie close to the break in the source counts. Two of the sources show narrow emission lines, and a third has a broad line. One of the narrow line-emitting sources has a clear redshift identification at \(z = 2.18\), while the other has a tentative determination based on the highest redshift detection of He I \(\lambda 10830\) at \(z = 1.26\). The remainder have featureless spectra to deep limiting equivalent widths of \(\sim 20–60\)\(\AA\) and line flux \(\sim 5 \times 10^{-17}\) erg s\(^{-1}\) cm\(^{-2}\) in the \(K\)-band. High-quality \(J\), \(H\) and \(K_s\)-band images of the sources were combined with archival optical detections or limits to estimate a photometric redshift for six. Two sources show complex double morphology. The hard sources have spectral count ratios consistent with heavily obscured AGN, while the host galaxy emits much of the optical and near-infrared flux. The most likely explanation for the featureless continua is that the line photons are being scattered or destroyed by optically-thick gas and associated dust with large covering fractions.

Key words: diffuse radiation – X-rays: galaxies – infrared: galaxies – galaxies: active

1 INTRODUCTION

The Chandra observatory has largely resolved the 2–10 keV X-ray background (XRB) within two years of its launch, after almost four decades of scientific effort (Giacconi et al. 1962; Mushotzky et al. 2000; Brandt et al. 2001; Giacconi et al. 2002). Follow-up work within the past year has revealed that this Chandra source population can be broadly split into type 1 active galactic nuclei (AGN), narrow emission-line AGN, optically-normal galaxies with no sign of activity other than in X-rays and optically-faint sources which are difficult to identify (Barger et al. 2001; Alexander et al. 2001; Rosati et al. 2001; Willott et al. 2001). XMM-Newton, with its higher effective area at 10 keV, has also begun to deliver results (Hasinger et al. 2001) which confirm the essential Chandra findings (e.g., Lehmann et al. 2001a; Mainieri et al. 2002).

\(^1\) Based on observations made with ESO Telescopes at the Paranal Observatories under programme 67.B-0188
\(^2\) Based on observations made with the Chandra X-ray Observatory

The individual X-ray spectra of the hard sources are flat enough to account for the XRB spectral slope of 0.4 (Marshall et al. 1980; Gruber et al. 1999), and their integrated flux contributes \(\sim 70 – 100\) percent of the XRB intensity. The ambiguity in the absolute intensity, still to be resolved, is caused by cross-calibration mismatches between various X-ray missions and/or cosmic variance of the sources themselves (e.g., Barcons et al. 2000; Cowie et al. 2002). Essentially, these observations agree with models which synthesize the XRB through the integrated emission of populations of obscured AGN that are spread over redshift and have a range of intrinsic absorbing column densities of gas (Setti & Walter 1989; Gilli, Salvati & Hasinger 2001; Wilman, Fabian & Nulsen 2000; Comastri et al. 1995). Little evidence had been found in pre-Chandra surveys for the existence of highly absorbed, intrinsically powerful objects (e.g., Halpern, Turner & George 1999).

1.1 Obscured ‘type 2’ sources

What would be the observable characteristics of such obscured objects? In the framework of the standard torus
We have been studying hard X-ray sources found serendipitously in the fields of massive clusters of galaxies with Chandra (Crawford et al. 2001; 2002; Gandhi et al. 2002). The flux regime that we sample is 10^{-12} - 10^{-11} cm^{-2}. This is an important regime, as the bulk of the XRB intensity is generated in populations with these fluxes (e.g., Cowie et al. 2002). Our work complements surveys such as ChaMP (Wilkes et al. 2001) and XMM-Newton serendipitous surveys (Watson et al. 2001; Barcons et al. 2001; Baldi et al. 2002).

By targeting the hard, optically-dim sources with multi-band imaging leading to photometric redshifts as well as optical spectroscopy, we have been able to discover both broad and narrow-line AGN and find that these lie at a large range of redshift (0.2 < z < 4). The broad-line sources have equivalent-widths and line-fluxes similar to those identified by previous missions such as ROSAT (Lehmann et al. 2001b) and are predominantly associated with the X-ray soft sources. The hard sources have X-ray spectra that suggest high but Compton-thin absorption and possess narrow emission lines (Seyfert 2s). We also find one source which does not fit this pattern: with high X-ray absorption (N_H > 10^{22} cm^{-2}) but a broad strong Mg II λ2797Å line, this source probably has a dust:gas ratio different to the Galactic value (Crawford et al. 2002).

The gravitational magnification of the cluster potential well enables us to study, rather than just detect, some sources which lie within 1 arcmin of the cluster centre. Two such sources magnified by factors of 2 and 8 by the cluster Abell 2390 were found to be powerful obscured (N_H > 10^{23} cm^{-2}) sources, even after de-magnification. The first source (A18 in Crawford et al. 2002) has an intrinsic X-ray luminosity L_{2-10 keV} > 10^{45} erg s^{-1}, while ISOCAM 6.7 and 15-μm detections are used to infer an absorbed big blue bump luminosity L_{UV} > 10^{45} erg s^{-1} for the second source (A15).

In this work we study a small sample of hard serendipitous sources in the near infrared (NIR) with the ISAAC instrument on the Very Large Telescope (VLT). There are a number of reasons for choosing the NIR regime: 1) as we have previously shown (Crawford et al. 2001), optically-dim sources are relatively bright in the NIR, making their study possible; 2) XRB synthesis models have so far predicted a characteristic source redshift of z = 2, where typically strong emission lines such as Hα or Hβ would be redshifted to the NIR (Fig 1); 3) longer wavelength rest-frame lines are less sensitive to reddening and should be detectable through higher obscuring columns.

This extends our previous work of obtaining UKIRT spectra (Crawford et al. 2001) of such sources to deeper equivalent width limits and provides constraints on the utility of using an 8m class telescope in the NIR for such studies.

Quoted cosmological quantities assume H_0 = 50 km s^{-1} Mpc^{-1} and q_0 = 0.5 throughout.

2 SAMPLE SELECTION

We are studying unresolved X-ray sources found serendipitously in Chandra ACIS-S observations of nearby galaxy clusters, preferentially selecting hard sources with count ratios ≤ 2 (see section 3). Roughly one-third of the entire ACIS sample is hard by this definition.

Our total sample consists of more than 20 clusters and covers the entire range in right ascension. Thus, our initial follow-up catalogue is the Digitized Sky Survey (DSS)^1, on which we target the sources that are dim or invisible (B>22.5; R>21; hereafter referred to as optically-dim). At the flux level that we sample, roughly 40 percent of all X-ray sources satisfy this criterion. As our pointings lie in the directions of massive galaxy clusters, we can make use of rich archival resources and often get fainter detections and/or limits. We find that optically-dim sources in 10-20 ksec ACIS exposures cover a wide range in magnitude: 22< B<24 and 19< I<23. Most of these are detected in the NIR at 16.5< K<19. For deeper Chandra exposures, our faintest limits are B>25.5, I>24 and K>20.

Our ‘best’ targets are the hard sources that are also optically-dim and we estimate these to be about 50 percent of the hard X-ray serendipitous sample. We also include a few exceptionally hard sources which may have optical counterparts on the DSS.

For the present work, we selected 8 sources found serendipitously in 3 clusters observable from the southern hemisphere (see next section). 6 of these are hard, and 2 soft sources are included for comparison. 6 are optically-dim and 1 soft source has R=20. We were forced to observe a relatively bright source with R = 19.6 (but very hard) due to poor atmospheric transparency on one night. This small

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1 http://archive.stsci.edu/cgi-bin/dss_form
sample represents about 13 percent of all the detected X-ray sources, and 25 percent of the hard sources.

### 3 X-RAY OBSERVATIONS

The Chandra cluster observations from which we chose sources for the present work are: MS 2137.3-2353 (sequence number 8000104, duration 50 ks), A 1835 (800003 of 19.6 ks) and A 2204 (800007, 10 ks). The effective exposure of the Chandra observation of MS 2137.3-2353 was limited by background flaring to 34.7 ks; the other observations were unaffected by background variations. The data were processed with the Chandra Interactive Analysis of Observations (CIAO\(^2\)) software, using versions 2.1 and 1.5. The Galactic line-of-sight column densities in the directions of MS 2137.3-2353, A 1835 and A 2204 are approximately 3.5, 2.3 and 5.6\(\times10^{20}\) cm\(^{-2}\) respectively (Stark et al 1992).

The source detection was performed in the total 0.5-7 keV band in a very similar manner to that described in Crawford et al (2002), using the standard WAVDETECT package in CIAO. WAVDETECT was run on the data in two forms: the original unbinned pixels (each of half-arcsec width), and binned by two pixels (ie arcsec-long pixels). We used the \(\sqrt{2}\) sequence of wavelet scales (ie 1, 1.414, 2.0 \ldots 16.0 pixels), and set the significance threshold for sources at \(10^{-6}\) (which implies that the expected number of false sources per ACIS chip is roughly one\(^3\)). We also checked that the 8 sources studied in this paper were detected with a lower threshold of \(10^{-7}\), thus making them unlikely to be spurious. We discarded all sources within 20 arcsec of the edge of each chip to avoid loss of source counts due to the spacecraft dither, and those with fewer than ten (non-background subtracted) counts. On the ACIS chips from which the present sample of sources was chosen, we found the following total number of sources in the respective cluster fields. MS 2137: 11 on S2[chip 6] and 25 on S3[7]; A 1835: 14 on S3[7] and 2 on I2[2]; A 2204: 8 on S2[6] and 3 on I3[3].

We also estimated the counts in each source in 3 energy bands: 0.5-2 keV (soft; S), 2-7 keV (hard; H) and 0.5-7 keV (total). The counts were taken from a box centred around a source, with a length given by the square root of the number of pixels in the source cell (as given from WAVDETECT). The local background was estimated from a concentric box with a length five times longer. The background box was

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\(^2\) http://cxc.harvard.edu/ciao/

\(^3\) WAVDETECT manual at http://cxc.harvard.edu/ciao/
The 38 counts have more both possibly interacting (Figure 4). The are unresolved. Is is A2204 and A2204 was also observed at 3, A1835 2 is relative and A1835 1, A1835 show clear on the subsequent nights for spectroscopy). We observed through the a small sample of hard sources (from which 4 were chosen since the field of view is 2.5 arcmin, only ratios that we observe (0.1-3; Table 1), the implied columns of the source, with the least counts being detected in the soft band where a $\Gamma = 2$ power-law model with only Galactic absorption predicts the most flux (dotted line).

The X-ray properties of the sources are listed in Table 1. None of these are magnified by strong lensing due to the cluster potential. Six of the sources are very hard (with S/H ratios of less than 1.6). Fig 2 shows how the S/H ratio can be translated into physical column densities under different assumptions. A simple absorbed power-law ($\Gamma = 2$) transmission model was adopted within XSPEC (Arnaud 1996) and counts in the soft and hard bands were predicted assuming different Galactic columns appropriate to the cluster line-of-sight and response matrices for each of the two kinds of CCDs on the Chandra ACIS instrument. For the ratios that we observe (0.1-3; Table 1), the implied columns are $5 \times 10^{22} \sim 10^{24}$ cm$^{-2}$ if at $z = 2$; or the redshifts are consistent with $0.5 < z < 2$ if $N_H = 5 \times 10^{22}$ cm$^{-2}$.

Only two sources (A 2204_1 and A 2204_2) have more than 30 counts in the total Chandra band. A2204_2 is relatively soft with most counts lying below 2 keV. A2204_1 is hard and in Fig 3, we show its extracted X-ray spectrum. Despite large errors, the spectrum clearly demonstrates the need for large amounts of X-ray absorption, as compared to the spectrum of a source emitting power-law (photon-index $\Gamma = 2$) radiation modified only by Galactic absorption appropriate to the line-of-sight to A 2204.

4 FOLLOW-UP IMAGING & PHOTOMETRY

4.1 Near-Infrared

Near-IR images of the selected sources in the MS 2137.3-2353 field were obtained using the imaging spectrograph ISAAC (Moorwood et al. 1997) on the Very Large Telescope (VLT). Since the field of view is 2.5 arcmin, only ~1 Chandra source can be imaged at a time. Thus, imaging was targeted toward a small sample of hard sources (from which 4 were chosen on the subsequent nights for spectroscopy). We observed through the J, H and Ks (hereafter K) filters on 2001 June 28 in seeing of ~0.3 arcsec (see Fig 4 for K-band thumbnail images of all the sources). The total integration time used was 600 s in H and K, and 720-900 s in J, with individual exposures of 10 s obtained in jitter mode around a grid with offsets of ~30 arcsec. Bad pixel map creation, dark current subtraction and flat-field division were carried out using the jitter routine of the eclipse software package V4.0.4 (Devillard 1997). In addition, jitter was used for background subtraction (using parameters suggested by Iovino 2001) and combination of jittered frames. Three photometric standard stars were observed over the course of the night and the zeropoint in each filter was found to be constant to within an RMS of 0.02 magnitudes.

Magnitudes were computed using the SExtractor package (Bertin & Arnouts 1996). We present the resulting magnitudes of our sources in Table 2 and a log of exposure times in Table 3. These are Kron magnitudes (with a Kron scale of 2.5; Kron 1980) for isolated sources and seeing-corrected isophotal magnitudes (with minimum isophote at 1.5r) for blended objects. It has been reported that the jitter pipeline may underestimate the brightness of sources in the K-band due to biasing of the local background in jittered images. We corrected for this by a simple prescription of increasing all the fluxes obtained at the VLT by 10 percent. Refer to Iovino (2001) for more details.

The excellent seeing conditions enabled us to clearly resolve some sources with a double morphology – MS2137_4 and A1835_1 – both possibly interacting (Figure 4). The signal-to-noise decreases in the H and J-bands, but the components are still resolved and separable for photometry. MS2137_3, A1835_1, A2204_1 and A2204_2 show clear ellipticity in the K-band, while the remaining sources – MS2137_1, 2 and A1835_2 are unresolved.

The sources in the field of A 2204 and A 1835 were imaged at UKIRT using the UFTI array (UT date 2000 Aug 11 and Feb 24 respectively; A1835_1 was also observed at ISAAC). A full description of the data reduction is given in Crawford et al. (2001). Essentially, this was carried out using the standard package CGS4DR V2.001 in a manner very similar to that for the ISAAC data. Magnitudes were measured in SExtractor, and are presented in Table 2.

4.2 Optical

DSS (both generation 1 and 2) blue and red images of the X-ray fields were downloaded directly from the ESO mirror of the DSS website. Deeper identifications of or detection limits to optical counterparts from the AAT, the CFHT and the INT were obtained from the respective archives. These data were calibrated using archival bias-subtraction and flat-fielding frames created on the night of observation. Fringing was removed from the I-band data by generating a master fringe frame from six offset frames for each of the archival datasets used.

Magnitudes were obtained using SExtractor as described in the previous section. Photometric standard stars were used for flux calibration (except one case; see Table 2). DSS magnitudes used a smooth extension of the first gener-

\[ \text{\url{http://archive.eso.org/dss/dss}} \]
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Figure 2. Predicted S/H ratios at various obscuring columns (Left, at $z = 2$) and redshifts (Right, at $N_H = 5 \times 10^{22} \text{ cm}^{-2}$) for an absorbed power-law transmission model ($\Gamma = 2$) and Galactic absorption appropriate to MS2137 (diamonds), A1835 (triangles) and A2204 (squares) as measured on the ACIS-S3 background-illuminated chip (solid) and the ACIS-S2 front-illuminated chip (dashed).

ation flux calibration by the Catalogs and Surveys Branch to fainter fluxes. Upper-limits for non-detected objects were defined as the flux corresponding to 3 times the background sky RMS in a 3-arcsec diameter aperture close to the source location. Such an aperture size is typical of the Kron apertures for the fainter of the detected sources. The final magnitudes are listed in Table 2. The INT Wide-Field Camera data was not fringe subtracted; thus apertures were selected by hand and background maps were inspected for consistency.

Typical seeing FWHM diameters for the optical observations were $\approx$1 arcsec. Combined with a plate scale which is typically larger than that of ISAAC, the morphology of the sources could not be reliably determined in the optical.

4.3 Source matching

All fields were cross-calibrated with the APM sky survey catalogue to generate an astrometric solution for each. RMS errors for the solutions were less than a pixel (better than the pixel scale) in all cases over large image regions.

Any astrometric offset between a Chandra source and its corresponding NIR counterpart is typically least ($\leq$ 1 arcsec) for sources on the ACIS-S3 chip 7, closest to the telescope aim-point. With off-axis PSF degradation, this may increase to several arcsec on the other chips. For example, the centroid determination of A1835-2 on ACIS chip 2 has large errors (Table 1), which leads to an offset of 5 arcsec from the source that we consider to be the NIR counterpart. There is, however, no source confusion problem (Fig 4) and identification is unambiguous (see also § 6.3.3). In all cases, we could identify a NIR counterpart and associated it with the nucleus of the object.

The probability of a false match occurring by chance was calculated by a ‘randomstep’ method similar to that used by Hornschemeier et al. (2001). Astrometric cross-correlation with DSS images was repeated after all X-ray sources were offset 10 arcsec to the north-east, north-west, south-west and south-east. The number of false source matches, averaged over the four offsets, was small: $\leq$ 0.5 (zero-offset source match numbers ranged between 10 and 20).

5 SPECTROSCOPIC OBSERVATIONS

We acquired spectra of the near-infrared counterparts to the Chandra X-ray sources during the nights of 2001 June 29 and 30, again using ISAAC on the VLT. We used a 1 x 120 arcsec slit and the low-resolution grating in the SK (hereafter K) and/or J bands. The choice of the low-resolution grating was made since no apriori information on the redshifts and detectability of the sources was known; the aim was to capture a large amount of source flux through the slit. This resulted in a dispersion of 7.138Å per pixel and a full-width at half-maximum (FWHM) of $\sim$45Å for a single narrow night-sky OH emission line assumed to be intrinsically unbroadened ($R \sim 450$) in K, and a dispersion of 3.610Å in J, with a limiting FWHM $\approx$25Å ($R \sim 500$). The first night had thick cloud at times, with seeing varying from 1 arcsec to greater than 3 arcsec, whereas the second night was photometric.

The objects were acquired by blind slit offsets from nearby bright stars as measured in pre-imaging. In many cases, the slit was long enough to encompass the star as well as target, thus providing a constant monitor of target acquisition. The targets themselves were nodded on the slit in an A–B–B–A pattern, with a small random jitter offset about each of two nod positions. Typical total exposure times of $\sim$2 hr were used, with each integration being 180 s long in order to be background limited. Given the significant loss of data quality on the first, non-photometric night, we only had time to obtain spectra in the J and K wavelength regimes, covering a wide choice of redshift space with two filters only.

Electrical ghost removal, dark current subtraction, distortion correction of sky lines, background subtraction, flat-
6 RESULTS

6.1 Colours and photometric redshifts

The hard S/H ratios argue against any of our targets being stars (The stellar locus is also well separated in \(J - H\) vs. \(H - K\) colours; c.f. a similar plot for sources in Crawford et al. 2002). The optical–infrared colours in Fig 7 show that most of the sources have colours or limits which are redder than the \(B - K\) colour of an unobscured quasar template (Elvis et al. 1994; median of a sample of UV-bright, radio-quiet quasars at low redshift). The red target colours are closer to those predicted for Coleman, Wu, Weedman (1980) redshifted templates (shown without reddening). The three sources with the bluest \(B - K\) limits (MS2137_1, MS2137_2 and A1835_1) are those with only relatively shallow optical lower-limits from the DSS.

We estimated photometric redshifts (\(z_{\text{phot}}\)) for all the sources using the publicly-available code HYPERZ (Bolzonella, Miralles & Pelló 2000), with the input parameters detailed in Crawford et al. (2002). Briefly, HYPERZ varies the redshift and reddening of various synthetic and empirical galaxy and AGN templates in order to best fit the observed fluxes. The synthetic templates used are Bruzual & Charlot (1993) models which are also evolved with redshift. The \(z_{\text{phot}}\) solutions are listed in Table 5, and shown in Fig 8.

If DSS upper-limits were shallow compared to the flux in the other filters, these were excluded from the HYPERZ fit (MS2137_4 and A1835_1). Note that 2 of the solutions (A2204_1 and A2204_2) are found at similar \(z_{\text{phot}} \approx 3.4\). This is primarily due to the large flux seen in the \(R\)-band relative to both \(B\) and \(J\), placing the Lyman break at \(z > 3\). No reliable redshift estimate could be obtained for the two sources (MS2137_1 and MS2137_2) with DSS upper-limits much shallower than the deep \(J, H\) and \(K\) magnitudes available.

6.2 Spectra

Three sources have strong detectable line emission features in their near-infrared spectra; the rest have no ambiguous
Figure 5. K-band ISAAC spectra. The y axis for all the target spectra begins at zero flux. The bottom two panels show the sky emission spectrum (at arbitrary scaling) and the sky transmission fraction (and hence the sky absorption). The sky emission spectrum has been effectively ‘flattened’ by dividing a polynomial to fit the rising thermal continuum in order to enhance the lines themselves. Lines expected at the derived redshift (spectroscopic or photometric) are marked in light grey.

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Figure 6. Same as in Fig 5 for sources observed in the J-band. The emission feature just before 1.15μm in all target spectra is most likely spurious (see text for details).
emission (or absorption) features above the sky noise (Figs 5, 6 and Table 6). Note that the apparent feature just red of $2\mu$m that shows up either as emission or absorption in many of the spectra (Fig 5) is not real. Its wavelength overlaps with that of the OH $\lambda 20005$ and $\lambda 20008$ 8–6 Q emission transitions, and also coincides with a strong telluric absorption feature. Since the telluric spectrum is effectively divided out during standard star division, a slight mismatch in the airmass between the target and the star, combined with variation in the intensity of sky OH emission on timescales of minutes, can produce such spurious features which are difficult to remove. Another such spurious emission feature occurs at $\approx1.145\mu$m in the $J$-band (Fig 6).

The limiting equivalent-widths (EW; observed frame) of the spectra range from 20–60Å in the $K$-band and are close to 15–20Å in the $J$-band, depending on the exposure time (except for MS2137.4 in the $J$-band which was observed for 1800 s only and the limiting EW is nearer 100Å). This
limit was estimated by measuring the equivalent width of an unresolved gaussian profile with amplitude approximately 3 times the standard deviation of the continuum at various positions along the dispersion axis. Note that this limit is shallower at the positions of sky emission lines.

A discussion of the three sources with unambiguous emission lines follows, while notes for the remaining objects/spectra can be found in the appendix.

6.3 Emission line sources

6.3.1 MS2137.1

MS2137.1 is invisible on the DSS images, but identified as a point source in the NIR. Its K-band spectrum (top spectrum in Fig 5; Fig 9 for detail) shows an unambiguous Hα+[NII] complex at 2.0846μm, implying a redshift of $z = 2.176±0.001$ for the source. Line emission from [SII]λ6717,6731 at the same redshift is marginally detected (at 2σ above the sky noise). While it appears that [OI]λ6300 is clearly detected, caution must be applied as this coincides with the spurious feature just beyond 2μm.

A model with discrete symmetric Gaussian profiles and identical velocity widths fixed to that of Hα overlaid on a constant continuum was fit using QDP (Tennant 1991). The Hα line is unresolved and substantially narrower than broad type 1 AGN permitted lines. The [NII]λ6584:Hα line intensity ratio is 0.56 and [SII]λ6717:Hα intensity is 0.12, which is consistent with that observed both in H II regions and Seyfert 2s (Veilleux & Osterbrock 1987; also Storchi-Bergmann 1991).

While observations of additional emission lines (e.g., redshifted [OIII]λ5007 in the H-band) may help to distinguish between starburst or AGN origin for the observed K-band emission lines, the hard X-rays must originate from a Seyfert 2. The implied X-ray (2–10 keV restframe) luminosity, assuming a $\Gamma = 1.4$ power-law at $z = 2.176$ affected only by Galactic absorption, is $L_{2-10} = 5.6 \times 10^{43}$ erg s$^{-1}$. If instead, we assume a crude model based on the observed S/H ratio of $\approx 1.1$, implying an obscuring column of $\approx 2 \times 10^{23}$ cm$^{-2}$ for $\Gamma = 2$ at $z \approx 2$ (Fig 2), the implied $L_{2-10} = 1.6 \times 10^{44}$ erg s$^{-1}$, after correcting for the obscuring column density.

6.3.2 A1835.1

A1835.1 is clearly resolved into two components on the ISAAC K-image (Fig 4). An astrometric solution associates the brighter westerly component (A1835.1W) with the X-ray source (though only marginally). The ISAAC spectrum of this brighter component has a single strong emission line at 2.4452μm, again unresolved (fifth spectrum from top in Fig 5). Although the night sky thermal emission begins to dominate beyond $\approx 2.35μm$, the narrow line is very significant (Fig 9) and lies fortuitously between two bright OH emission lines. The line connects both components of the source and is thus spatially extended (we were able to obtain the spectra of both component simultaneously with a slit orientation of 90°).

We attempted to identify weak features in the spectrum without much success, before trying to use the photometric redshift estimated from HYPERZ as a first guide towards possible line identification. $z_{\text{phot}} = 1.2\pm0.15$ (Table 5) suggests that the strong line at 2.4452μm and a weak $\sim 3σ$ emission feature at 2.2673μm could be He I λ10830 and Paα λ10049 respectively (Fig 9), implying a tentative redshift $z = 1.256$, well within the 90 percent $z_{\text{phot}}$ confidence interval. The He I transition has been observed previously in Seyfert 2s with intensity approaching Hβ (and even exceeding it in NGC 1068 and Mrk 3; Rudy et al. 1989; Osterbrock, Shaw & Veilleux 1990). Its strength is enhanced due to collisional excitation from the metastable 2$^3P$ level to 2$^3P$, followed by radiative decay (Osterbrock 1989).

The other strong feature expected in this regime ([SIII]λ9531) is not observed by us. This may imply that the density of the narrow-line region is much greater than the critical density to collisional de-excitation $N_{\text{critical}} \approx 6 \times 10^6$ cm$^{-3}$. If our identification is confirmed, this would be the highest redshift observation of He I λ10830 to date.

Estimating the intrinsic luminosity by correcting for a column-density of $\approx 10^{23}$ cm$^{-2}$ as implied by the S/H ratio of 0.41 (on a front-illuminated chip) at the inferred redshift, we deduce the presence of a powerful Seyfert emitting power-law radiation ($\Gamma = 2$ fixed) with $L_{2-10} = 1.2 \times 10^{44}$ erg s$^{-1}$.

6.3.3 A1835.2

The large $\sim 5$ arcsec offset between the Chandra source A1835.2 and its counterpart is most likely due to the difficulty of centroiding with off-axis ($\approx 12$ arcmin) PSF degradation. The formal probability of a match at such a large offset is 1%. However, if this were not the counterpart, a large X-ray/optical flux ratio and corresponding optical/NIR extinction would be inferred since the counterpart would be invisible in our images. This is unlikely given the softness of the X-ray source.

The spectrum (sixth from top in Fig 5) – is not easy to interpret. There is a clear broad (FWHM$\approx 4600$ km s$^{-1}$)
feature centred at 2.3480\,\mu m and two ambiguous features marked with a '?' in the figure. No obvious emission line pattern fits these data, and it is likely that the ambiguous lines are spurious and caused by telluric absorption in the first case, and increased thermal noise in the second.

The presence of a broad feature (Fig 9 for detail), the high S/H ratio (Table 1) and the point-like morphology (Fig 4) suggests that we are viewing a quasar. Fitting a QSO template to the broad-band colours gives an acceptable fit at $z_{\text{phot}}=3.4$ [3.0, 4.0] (Table 5). Thus a potential match for the broad feature is with H$\beta$ for $z=3.830 \pm 0.005$ (within the 90 percent $z_{\text{phot}}$ interval). Any [OII]35007 emission would overlap with an OH emission feature close to 2.1488\,\mu m (refer to Fig 5; if this [OIII] line exists, it must have EW<40\,\AA in the observed frame). We do not see any H$\gamma$A4340 emission. Its predicted photoionization recombination intensity is 0.47 times that of H$\beta$, though we note that Vanden Berk et al. (2001) found that EW(H$\gamma$) $\approx$ 0.25$\times$EW(H$\beta$), which implies an H$\gamma$ strength close to our limiting EW (Table 6). Moreover, any reddening is likely to further decrease the H$\gamma$:H$\beta$ intensity ratio. Thus, its absence may not be surprising.

Again, assuming $\Gamma = 2$ and correcting for a column-density of $\approx 7 \times 10^{22}$ cm$^{-2}$ (S/H = 3.1), the implied $I_{\gamma-10} \approx 6 \times 10^{44}$ erg s$^{-1}$ at $z_{\text{phot}}=3.830$, which is consistent with quasar luminosity.

7 DISCUSSION

Of the 8 objects, we detect a broad K-band emission line in one source (A1835_2) and narrow lines in two sources (MS2137_1 and A1835_1). The spectrum of the narrow-line source MS2137_1 places a firm redshift constraint of $z = 2.18$. The other two sources have only a single unambiguous line in their spectra and tentative identifications (at $z = 1.26$ and 3.8) for these lines are determined in conjunction with the photometric redshifts. Two sources (MS2137_3 and A2204_1) have high signal/noise $K$ and K-band featureless continua. The limiting observed-frame equivalent-widths (EW) for these are $\sim 20\,\AA$ in both bands. For the remaining sources, the limiting EW range over $\sim 30$-$60\,\AA$ in the K-band. Few useful constraints can be obtained from the noisy spectra of MS2137_2 and MS2137_4 (J-band).

The high quality NIR images reveal that the sources have a mixture of unresolved (MS2137_1, MS2137_2 and A1835_2), galaxy-like (MS2137_3, A2204_1 and A2204_2) and double (MS2137_4 and A1835_1) morphologies.

To investigate the lack of significant spectral features in some sources, we ask the question: how likely is it that an emission line from a Seyfert galaxy at a randomly chosen redshift would be redshifted into the $J$ or $K$ bandpasses? Fig 1 goes some way in answering this question. Most of the transitions mentioned in the figure are strong lines in Seyfert spectra, and all can easily have an EW greater than our best limiting value of 20\,\AA. There are few regions in redshift space where no redshifted lines are observed in either (or both) of the ISAAC bandpasses. Especially important is the redshift regime $z \lesssim 1$, as suggested by recent Chandra and XMM-Newton findings (e.g., Rosati et al. 2001, Brandt et al. 2002, Hasinger 2002). If the sources are at $z \lesssim 1$, we would see lines ranging from H$\alpha$$\lambda$6563 to P$\alpha$$\lambda$1875. Table 7 lists these lines and their typical strengths. H$\alpha$ and [NII] are the strongest expected lines, followed by He I and forbidden transitions of Sulphur.

In summary, at $0.7 < z < 1$, we would observe H$\alpha$+[NII]. Detection of other (typically weaker and/or lower $z$) lines is also likely, as they have been observed in the literature with large strengths and EW. Thus, although we cannot rule out the possibility that we miss the lines at $z < 1$, we consider alternative possibilities as well.

Early-type galaxies with little on-going star-formation show many absorption features in their spectra (that we would not detect due to signal/noise constraints) but few (if any) emission features. This is especially true for extremely-red objects (EROs). In a sample of ISAAC spectra of galaxies selected for their extremely red colours, Cimatti et al. (1999) found neither strong emission lines nor continuum breaks. While it may be the case that our NIR spectra are dominated by galactic light, the hard X-ray luminosities (assuming redshifts as determined from photometric or spectroscopic constraints) range over $10^{39} - 6 \times 10^{44}$ erg s$^{-1}$. The X-rays must, therefore, originate in an active Seyfert-like nucleus, and the absence of any optical/NIR lines associated with the AGN must be accounted for. We consider the possibility that large column density gas and associated dust scatters / absorbs these line photons.

As discussed in section 3, the hard S/H ratios (Table 1) suggest high obscuration of the nucleus. Intrinsic columns below $\sim 10^{22}$ cm$^{-2}$ will result in S/H ratios which are
softer than observed ($\geq 3.5$) for all $z \geq 1$, if $\Gamma = 2$. The 6 hard sources must have $N_{H} > 5 \times 10^{22}$ cm$^{-2}$ if at $z = 2$ or $N_{H} > 10^{22}$ cm$^{-2}$ ($A_{V} > 6$ mags assuming a Galactic dust:gas ratio) if at $z = 0.5$. We note that evidence for such a dusty environment was found for at least one type 2 QSO in the field of A2390 by Wilman, Fabian & Gandhi (2000) through radiative transfer modelling based on photometric detections in the ISOPHOT mid-IR bands (cf. Crawford et al. 2002). The column density of the gas as inferred by the optical depth of the dust was found to be consistent with the X-ray measurements. A model incorporating dust in narrow-line region clouds was proposed by Netzer & Laor (1993).

The very high ratio of 0.16 for A2204.1 (a source with high signal:noise featureless continuum) suggests a column $> 10^{23}$ cm$^{-2}$ for all $z \geq 0.5$. In fact, extrapolating the de-absorbed flux at 1 keV for a $\Gamma = 2$ power-law model to $B$-band flux assuming a ‘typical’ quasar broad-band energy distribution (i.e. $\alpha_{B \lambda}$ spectral index; Elvis et al. 1994) implies $B_{\text{predicted}} \approx 19.5$ if A2204.1 lies at $z = 0$, or $B \approx 20.2$, if at $z = 2$. For the observed DSS limit of 22.5, this implies at least 2.5 optical magnitudes of extinction to the nucleus.

Compact, very optically-thick obscuration of the kind proposed by Pier & Krolik (1992) will obscure emission from the nucleus itself. If obscuration is in the form of moderately thick tori spread over tens or hundreds of parsecs (Granato, Danese & Franceschini 1994), line emission from star-forming regions could be partially absorbed as well. However, the distinct lack of emission lines (in both the optical and NIR), in these and other recently-found hard X-ray sources, requires covering fractions that are close to 4%. As opposed to low-redshift Seyferts, obscuration at high redshift would then be mostly independent of orientation. Indeed, large covering fractions (85 percent) have been predicted by Fabian & Iwasawa (1999), by correcting for absorption in the spectrum of the XRB. Forthcoming observations in the far-IR will significantly improve our understanding of the scale and extent of the enshrouding gas/dust distribution (e.g., SIRT; Brandl et al. 2000).

8 CONCLUSIONS

We have obtained VLT ISAAC spectra of 8 Chandra X-ray sources and detect continuum emission from all. These include 7 spectra in the K-band and 4 in J. 6 of these are optically-dim and the remaining 2 are dim in the $B$-band. The X-ray spectral count ratios constrain the obscuring column; 6 sources are very hard with the hardest being consistent with an intrinsic obscuring column density $N_{H} \geq 10^{23}$ cm$^{-2}$. These are typical of the population which contributes the maximum flux per source to the X-ray background.

We have been able to identify 2 narrow-line AGN (one at $z = 2.18$ and one possibly at $z = 1.26$; both consistent with luminous Seyfert 2s) and one broad-line AGN possibly at $z = 3.83$ for which intrinsic quasar luminosity is inferred. This includes possible detection of the most distant He I $\lambda 10830$ emission to date. Spectra in other wavebands will help to confirm these. Photometric redshifts have also been determined.

Even in the above long integrations ($\sim 1$ or 2 hours long) on an 8-m telescope, we are able to detect significant emission lines in only 3 of the 8 sources and identify the redshift of one of these unambiguously. This extends our previous 4-m UKIRT observations to better equivalent-width limits and confirms our earlier findings that although such sources are readily observed in the near-infrared, detailed identification of the source type is not straightforward. The absence of emission lines in many type 2 AGN, which contribute a large fraction of the XRB intensity, is evidence for high covering fractions of intrinsic obscuring gas.

9 ACKNOWLEDGEMENTS

The work presented here is based on observations obtained with the Chandra telescope and the VLT. In addition, use is made of archival imaging data from a number of telescopes (the AAT, INT, WHT and CFHT) and the DSS. It is a pleasure to acknowledge the support received from all these organizations in the process of acquiring our data. The eclipse team at ESO is thanked for constant help with ISAAC data reduction. We are grateful to the referee for thorough and constructive criticism. PG would like to thank Andrew Bunker for useful discussions and the Isaac Newton Trust and the Overseas Research Trust for support. CSC and ACF acknowledge financial support from the Royal Society.

APPENDIX: NOTES ON SOURCES WITHOUT SIGNIFICANT EMISSION LINES

**MS2137 2**

By an appropriate orientation of the long slit, MS2137.2 (second spectrum from top in Fig 5) was observed simultaneously with MS2137.1. Unfortunately, the only optical photometric data available to us is the DSS, on which, like MS2137.1, the source is invisible. Compared to the deeper ISAAC $J$, $H$ and $K$ detections of this source, the $B$ and the $R$ upper-limits are too shallow to obtain a satisfactory photometric redshift solution. The source is very dim and point-like in $K$, with no obvious spectral features down to a 3$\sigma$ limit on the EW of $\sim 40\AA$.

**MS2137 3**

MS2137.3 (third K spectrum from top in Fig 5 and top spectrum in Fig 6) is the only source to show an obvious slope in the continuum, which rises gently toward the blue in the K-band and drops suddenly at $\sim 2\mu$m. Though reminiscent of a P-Cygni profile, the proximity of the ‘break’ to the OH emission at 2.0008$\mu$m as well as the deep telluric absorption strongly suggests that this drop is due to the sky. The signal:noise in the continuum is $\sim 10$ per pixel, which may also explain the lack of significant absorption features. (The typical depth of such features expected at $z \approx z_{\text{phot}}$ in the K-band is also about 10 percent of the continuum; e.g., the temple NIR galaxy spectrum of Mannucci et al. 2001)

**MS2137 4**

MS2137.4 lies in an apparently interacting system (Fig 4) of galaxies with very red $B - K$ colour and diffuse tidal
structure in the K-band image. Astrometric calibration associates the X-ray source with the fainter, more north-easterly component. The signal/noise in the continuum spectrum (fourth from top in Fig 5) is comparable to that of MS2137_1; however, unlike MS2137_1, we do not detect any emission features, though HJ would lie in the K-band, if \( z = z_{\text{phot}} = 3.3 \). The source is then red in \( B - K \) due to the \( k \)-correction of the Lyman break (Fig 8). Due to time constraints, only a 30-minute exposure could be obtained in the J-band (fourth spectrum from top in Fig 6). This is also the faintest J-band source observed spectroscopically, which implies a poor equivalent width lower-limit \( EW_{\text{lim}}^{3.4} \approx 100\AA \) for any real features.

**A 2204_1 and A 2204_2**

A 2204_1 has a flat, featureless spectrum in both the \( J \) and \( K \) bands (seventh spectrum from top in Fig 5 and third spectrum from top in Fig 6), despite being the brightest NIR source in our sample (It is also the second brightest and the hardest X-ray source; Fig 3 for the X-ray spectrum). Unfortunately, the ISAAC spectra of this source were observed on a non-photometric part of the night with seeing \( \approx \) 3 arcsec, which led to severe degradation of the signal/noise, thus precluding possible identification of absorption features (The 1\( \sigma \) photon noise variation is \( \approx \) 8 percent of the continuum in the best part of the spectrum close to 2.2\( \mu m \)). Intriguingly, the blue \( R - J = 1.3 \) colour predicts a relatively young system (Fig 8), though we see no emission lines due to young stars. At \( z = z_{\text{phot}} = 3.4 \) (Table 2), H\( \beta \) and the \([\text{OIII}]\lambda\lambda 4959,5007 \) doublet are expected to fall in the K-band and Mg II \( \lambda 2797 \) in the J-band. Note that the K-band morphology clearly indicates an extended source. At \( z = z_{\text{phot}} \), a half-light diameter \( \approx 7 \) kpc is implied, which is large; but such large values have been reported in the literature (e.g., Ivison et al. 2001).

A 2204_2 is a soft source (Table 1). It is also at \( z_{\text{phot}} = 3.4 \), since its colours and red \( B - R \) limit are similar to A2204_1. A single J-band spectrum only 1-hour long was obtained due to time-constraints and no significant (non-spurious) features are found down to an \( EW_{3.4}^{\text{lim}} \approx 17\AA \) (fourth spectrum from top in Fig 6).

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Crawford C.S., Fabian A.C., Gandhi P., Wilman R.J.,


Crawford C.S., Gandhi P., Fabian A.C., Wilman R.J.,


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Moorwood A. et al., 1997, SPIE, 2871, 1146
Rosati P. et al., 2002, 566, 667
Table 1. Cluster/observation

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<th>RA</th>
<th>DEC</th>
<th>DEC</th>
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<th>Hard</th>
<th>S/H</th>
<th>Distance (arcsec)</th>
<th>0.57 keV Flux</th>
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<td>(2)</td>
<td>(3)</td>
<td>(4)</td>
<td>(5)</td>
<td>(6)</td>
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<td>(12)</td>
</tr>
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<td>21 40 29.3</td>
<td>(0.7)</td>
<td>23 47 41</td>
<td>(0.0)</td>
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<td>(1.0)</td>
<td>23 47 46</td>
<td>(0.7)</td>
<td>4.15</td>
<td>13.30</td>
<td>(4.4)</td>
<td>8.17</td>
<td>(3.19)</td>
<td>5.21</td>
<td>(3.11)</td>
</tr>
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<td>21 40 22.2</td>
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<td>21 35 03</td>
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<td>12.73</td>
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<td>(3.82)</td>
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<td>4.06</td>
<td>18.81</td>
<td>(7.49)</td>
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<td>(6.25)</td>
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<td>14.40</td>
<td>38.60</td>
<td>(6.30)</td>
<td>5.17</td>
<td>(2.54)</td>
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<td>(5.81)</td>
</tr>
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<td>A220.1</td>
<td>6</td>
<td>16 31 58.6</td>
<td>(0.9)</td>
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<td>6.00</td>
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<td>(7.16)</td>
<td>35.17</td>
<td>(6.13)</td>
<td>11.49</td>
<td>(3.67)</td>
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</table>

Column 7 denotes counts in the total 0.5-7 keV band.
Column 13 gives the ratio of the counts in the soft band to that in the hard band.
The fluxes in column 16 were derived using TIMMS assuming that each source has a power-law spectrum with slope of $\Gamma = 1.4$ with absorption only due to the Galactic line-of-sight column density.

The X-ray coordinates of this source reported in Crawford et al., 2001 are erroneously placed 9 arcsec away.

Table 2. Photometry (Vega magnitudes)

<table>
<thead>
<tr>
<th>Object</th>
<th>B</th>
<th>R</th>
<th>I</th>
<th>J</th>
<th>H</th>
<th>K</th>
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<td>MS217.7</td>
<td>24.5±0.2</td>
<td>0.45</td>
<td>22.1±0.2</td>
<td>21.9±0.3</td>
<td>20.4±0.3</td>
<td>22.1±0.2</td>
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<td>MS217.7</td>
<td>&gt;24.5</td>
<td>&gt;21.0</td>
<td>&gt;22.1</td>
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<tr>
<td>A1835.1</td>
<td>&gt;22.5</td>
<td>24.0±0.1</td>
<td>&gt;22.5</td>
<td>&gt;22.5</td>
<td>&gt;22.5</td>
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<td>&gt;21.0</td>
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<td>&gt;22.5</td>
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<td>A220.1</td>
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<td>&gt;22.5</td>
<td>&gt;22.5</td>
<td>&gt;22.5</td>
</tr>
</tbody>
</table>

Magnitudes marked by the following superscripts are obtained from data from the following telescopes (see Table 4 for details)

A: AT
B: CFHT
C: DSS
D: INT Prime Focus
E: UKIRT
F: VLT
G: VLT ISAC
H: INT WFC

1: Magnitude is based on a single flux standard observation, cannot resolve the two components.
2: Aperture magnitude
3: Photometric standard star unavailable, Standard zero points assumed, since observing conditions were reported to be photometric by the Carnegie Meridian Telescope Extinction Monitor [http://www.jwst.com/uk/3d/meridian/]

Crawford et al., 2001a
Table 3. Near-IR Observation Log

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<tbody>
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<td></td>
<td>J (s)</td>
<td>K (s)</td>
<td>Slit PA (°)</td>
<td>J (s)</td>
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<td>31</td>
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The coordinates denote the position measured in the near-IR pre-imaging, which was carried out at the VLT for objects in the field of MS2137, and at UKIRT for the rest of the sources.
The slit position angles are given in degrees east of North.

Table 4. Photometric Observations

<table>
<thead>
<tr>
<th>Band</th>
<th>UT Date</th>
<th>Telescope &amp; Instrument</th>
<th>Plate Scale (arcsec/pixel)</th>
<th>Filters</th>
<th>Seeing (arcsec)</th>
<th>Typical Exposure (seconds)</th>
<th>m_{lim}</th>
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<td>B</td>
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<td>AAT Prime Focus</td>
<td>0.391</td>
<td>KPNO 1</td>
<td>1.3</td>
<td>600</td>
<td>25.4</td>
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<tr>
<td>R</td>
<td>1994 Jun 07</td>
<td>INT Prime Focus</td>
<td>0.590</td>
<td>Kitt Peak 3</td>
<td>2.8</td>
<td>600</td>
<td>20.7</td>
</tr>
<tr>
<td>R</td>
<td>1998 Feb 26</td>
<td>CFHT STIS2</td>
<td>0.439</td>
<td>CFHT #4609</td>
<td>1.1</td>
<td>600</td>
<td>25.5</td>
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<tr>
<td>I</td>
<td>2000 May 01</td>
<td>INT WFC Prime</td>
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<td>Sloan i</td>
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<td>1198</td>
<td>22.9</td>
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<td>J, H, K</td>
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<td>0.6</td>
<td>540</td>
<td>21.4, 20.7, 20.0</td>
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<tr>
<td>J, H, K</td>
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<td>VLT ISAAC</td>
<td>0.148</td>
<td>J, H, Ks</td>
<td>0.3</td>
<td>600</td>
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The last column states the limiting magnitude as defined by 3 times the background sky RMS in a 3 arcsec aperture.

Table 5. Photometric Redshifts

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<th>\chi^2 (filters)</th>
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<th>A_V</th>
<th>M_B</th>
<th>z_{secondary}</th>
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<td>2.5 (6)</td>
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<td>0.7</td>
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<td>2.7 (5)</td>
<td>Burst</td>
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<tr>
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<td>1.23 (1.0, 1.5)</td>
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<td>0.2</td>
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<td>3.40 (3.0, 4.0)</td>
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<td>-</td>
<td>0.6</td>
<td>-25.2</td>
<td>3.6</td>
</tr>
<tr>
<td>A2204_1</td>
<td>3.40 (3.3, 3.6)</td>
<td>0.8 (5)</td>
<td>Burst</td>
<td>0.1</td>
<td>0.0</td>
<td>-27.3</td>
<td>0.5</td>
</tr>
<tr>
<td>A2204_2</td>
<td>3.39 (2.9, 4.1)</td>
<td>1.2 (5)</td>
<td>Burst</td>
<td>0.1</td>
<td>0.0</td>
<td>-26.4</td>
<td>0.5</td>
</tr>
</tbody>
</table>

The 90 percent \(z_{phot}\) confidence interval in parentheses in column 2 assumes \(\Delta \chi^2 = 2.7\).
The ‘filters’ parentheses in column 3 states the number of photometric data points used in the fit.
M_B in column 8 is the absolute Vega Magnitude in the B Bessell filter.
The Burst template is a single burst of star-formation at zero age followed by passive evolution.
The QSO template is from Francis et al. (1991).
\(^{\dagger}\)With the inclusion of the R-flux, HYPERZ prefers this solution to the one inferred by Crawford et al. (2001; based on 4 filters).

Table 6. Observed Spectroscopic Features

<table>
<thead>
<tr>
<th>Object</th>
<th>Feature</th>
<th>Width (σ)</th>
<th>FWHM</th>
<th>EW</th>
<th>Flux</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td>A2.0846</td>
<td>18.5</td>
<td>12.8</td>
<td>Ho 6563</td>
<td>= &gt; z = 2.176</td>
<td></td>
<td></td>
</tr>
<tr>
<td>A2.0916</td>
<td>18.5(^{\dagger})</td>
<td>7.2</td>
<td>NI 6554</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>A2.1355</td>
<td>18.5</td>
<td>3.9</td>
<td>[SII] λ6717 + 6731</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>A2.4452</td>
<td>18.1</td>
<td>8.7</td>
<td>He I 10830</td>
<td>= &gt; z = 1.256</td>
<td></td>
<td></td>
</tr>
<tr>
<td>A2.2673</td>
<td>20.9</td>
<td>1.9</td>
<td>H I Pa 7 λ10049</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>A2.3480</td>
<td>152.5</td>
<td>36.9</td>
<td>H3α4861</td>
<td>= &gt; z = 3.830</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\(^{\dagger}\)Velocity width fixed to that of Ho.
The numbers in square brackets give the 90 percent confidence interval for a single parameter.
Strengths stated in row 3 are for the combined [SII] doublet.
The line width and equivalent-width in the rest-frame will be lower than the above values by a factor of \((1 + z)\).
Table 7. Detectability of emission lines at $z \leq 1$

<table>
<thead>
<tr>
<th>Transition</th>
<th>J</th>
<th>K</th>
<th>$I(\text{transition})/I(H\beta)$</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Hαλ6563</td>
<td>0.71</td>
<td>&lt; 1.06</td>
<td>2.87</td>
<td>(1)</td>
</tr>
<tr>
<td>[NII]λ6584</td>
<td>0.70</td>
<td>&lt; 1.05</td>
<td>≈2.8</td>
<td>(2), (1)</td>
</tr>
<tr>
<td>[SII]λ6731</td>
<td>0.66</td>
<td>&lt; 1.01</td>
<td>0.5</td>
<td>Mrk 1073; (3)</td>
</tr>
<tr>
<td>[SIII]λ9069</td>
<td>0.24</td>
<td>&lt; 0.49</td>
<td>1.18 &lt; $z &lt; 1.71$</td>
<td>0.1</td>
</tr>
<tr>
<td>[SIII]λ9531</td>
<td>0.18</td>
<td>&lt; 0.42</td>
<td>1.07 &lt; $z &lt; 1.58$</td>
<td>0.25</td>
</tr>
<tr>
<td>He I λ18020</td>
<td>0.03</td>
<td>&lt; 0.25</td>
<td>0.83 &lt; $z &lt; 1.27$</td>
<td>0.78</td>
</tr>
<tr>
<td>Paβλ12822</td>
<td>–</td>
<td></td>
<td>0.54 &lt; $z &lt; 0.92$</td>
<td>0.17</td>
</tr>
<tr>
<td>Paαλ18756</td>
<td>–</td>
<td></td>
<td>$z &lt; 0.31$</td>
<td>0.35</td>
</tr>
</tbody>
</table>

Column 4 denotes the intensity with respect to that of Hβ. Any observed reddening has been corrected for. The last column gives the reference for the intensity calculation and/or measurement as follows:

(1) Case B photoionization prediction
(2) Storchi-Bergmann 1991
(3) Rudy et al. 1989
(4) Osterbrock & Veilleux 1989
(5) Goodrich, Veilleux & Hill 1994